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## Divergence formulas (conservation laws) for families of curves and surfaces and applications<sup>\*</sup>

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Abstract. The divergence formulas we have obtained (differential conservation laws) of the form div  $\mathbf{F} = 0$  for an arbitrary smooth field of unit vectors  $\boldsymbol{\tau}(x, y, z)$ , for a family of spatial curves as well as  $\{L_{\tau}\}$  for a family of surfaces  $\{S_{\tau}\}$  continuously filling a certain domain. The solenoidal vector field  $\mathbf{F}$  in these formulas is expressed, respectively, through the field  $\boldsymbol{\tau}(x, y, z)$ , the characteristics of the curves  $L_{\tau}$  and characteristics of the surfaces  $S_{\tau}$ . Also, we have found the formulas connecting the surface characteristics and those of the curves orthogonal to them. In the case when the curves  $L_{\tau}$  and the surfaces  $S_{\tau}$  are vector lines of the vector field  $\boldsymbol{v} = |\boldsymbol{v}|\boldsymbol{\tau}$  with the direction  $\boldsymbol{\tau}$  and the surfaces orthogonal to them, the conservation laws found are equivalent to divergence formulas for the field  $\boldsymbol{v}$ . With these general geometric formulas the divergent identities (differential conservation laws) for the solutions of the eikonal equation  $|\operatorname{grad} \tau|^2 = n^2(x, y, z)$ , the Poisson equation  $u_{xx} + u_{yy} + u_{zz} = -4\pi\rho(x, y, z)$  and for solutions of Euler's hydrodynamic equations are obtained. In the plane case, these formulas transform to the conservation laws obtained earlier.

This paper is a generalization and development of the published works [1–4].

The vector line  $L_{\tau}$  of vector fields corresponding to the solutions of the mathematical physics equations, and to the surfaces  $S_{\tau}$  orthogonal to them with the normal  $\tau$  continuously fill the domain in question. Therefore, in this paper we study not only the properties of fixed curves and surfaces, but the properties of their families.

In [4], we obtained divergence formulas (conservation laws) for the family  $\{L_{\tau}\}$  of the plane curves  $L_{\tau}$  with the tangent unit vector  $\boldsymbol{\tau} = \boldsymbol{\tau}(x, y)$  and the unit normal  $\boldsymbol{\nu} = \boldsymbol{\nu}(x, y)$  or for an arbitrary smooth field of the unit vectors  $\boldsymbol{\tau}(x, y)$  with the vector lines  $L_{\tau}$ . These conservation laws have the form div  $\boldsymbol{S}(\boldsymbol{\tau}) = 0 \Leftrightarrow \operatorname{div} \boldsymbol{S}^* = 0$ , where  $\boldsymbol{S}(\boldsymbol{\tau}) \stackrel{\text{def}}{=} \operatorname{rot} \boldsymbol{\tau} \times \boldsymbol{\tau} - \boldsymbol{\tau} \operatorname{div} \boldsymbol{\tau} = \boldsymbol{S}^* \stackrel{\text{def}}{=} \boldsymbol{K}_{\tau} + \boldsymbol{K}_{\nu}, \, \boldsymbol{K}_{\tau} = (\boldsymbol{\tau} \cdot \nabla)\boldsymbol{\tau} = \operatorname{rot} \boldsymbol{\tau} \times \boldsymbol{\tau} = k\boldsymbol{\nu}, \, \boldsymbol{K}_{\nu} = (\boldsymbol{\nu} \cdot \nabla)\boldsymbol{\nu} = \operatorname{rot} \boldsymbol{\nu} \times \boldsymbol{\nu} = -k_{\nu}\boldsymbol{\tau}$  are curvature vectors of the curves  $L_{\tau}$  with the curvature k and the curves  $L_{\nu}$  orthogonal to them with the tangent unit vector  $\boldsymbol{\nu}$  and the vector products of the vectors  $\boldsymbol{a}$  and  $\boldsymbol{b}, \nabla$  is the Hamiltonian operator,  $(\boldsymbol{v} \cdot \nabla)\boldsymbol{a}$  is the derivative of the vector  $\boldsymbol{a}$  in the direction of the unit vector  $\boldsymbol{v}$ .

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In this paper, we discuss the three-dimensional case, when we have the field of unit vectors  $\boldsymbol{\tau} = \boldsymbol{\tau}(x, y, z)$ , the family of spatial curves  $L_{\tau}$  with the Frenet basis  $(\tau, \nu, \beta)$  [5–7]  $(\tau$  is the unit tangent vector,  $\nu$  is a principal normal,  $\beta$  is binormal), the first curvature k and the second curvature  $\varkappa$ , and the family  $\{S_{\tau}\}$  of the surfaces  $S_{\tau}$ , orthogonal to the curves  $L_{\tau}$ , with the normal  $\boldsymbol{\tau}$ , the principal directions  $\boldsymbol{l}_1$  and  $\boldsymbol{l}_2$ , the principal curvatures  $k_1$ and  $k_2$ , the mean curvature  $H \stackrel{\text{def}}{=} (k_1 + k_2)/2$  and the Gaussian curvature  $K \stackrel{\text{def}}{=} k_1 k_2$  [5–7]. All the values  $\boldsymbol{\tau}, \, \boldsymbol{\nu}, \, \boldsymbol{\beta}, \, k, \, \varkappa$  and  $\boldsymbol{l}_1, \, \boldsymbol{l}_2, \, k_1, \, k_2, \, H, \, K$  are vector and scalar fields in D, which is continuously filled by the curves  $L_{\tau}$ and the surface  $S_{\tau}$ . In the three-dimensional case the value of div  $S(\tau)$  is not identically equal to zero in D. The following geometric reason explains this circumstance. In [8], it was found that the value of  $\{-\operatorname{div} S(\tau)/2\}$  is the Gaussian curvature K of the surface  $S_{\tau}$ . The plane case corresponds to the cylindrical surfaces  $S_{\tau}$  with the directrices  $L_{\nu}$  and generatrices parallel to the axis Oz; their Gaussian curvature  $K \equiv 0$  and hence, div  $S(\tau) = 0$ in D. However, in the general case,  $K \neq 0$  for a single surface  $S_{\tau}$  (except for evolving surfaces [5–7]) and, especially, for the family  $\{S_{\tau}\}$ , i.e., in D; hence, div  $S(\tau) \neq 0$  in D.

In this paper, we obtain the divergence formula (differential conservation laws) of the form div F = 0 for an arbitrary smooth field of the unit vectors  $\tau(x, y, z)$ , for the family of spatial curves  $\{L_{\tau}\}$  and for the family of the surfaces  $\{S_{\tau}\}$  with the normal  $\boldsymbol{\tau} = \boldsymbol{\tau}(x, y, z)$ . The solenoidal vector field  $\boldsymbol{F}$ in these formulas is expressed, respectively, through the field  $\tau(x, y, z)$ , the characteristics  $\boldsymbol{\tau}, \boldsymbol{\nu}, \boldsymbol{\beta}, \boldsymbol{k}, \boldsymbol{\varkappa}$  of the curves  $L_{\tau}$  and the characteristics  $\boldsymbol{l}_1, \boldsymbol{l}_2$ ,  $k_1, k_2, K, H$  of the surfaces  $S_{\tau}$ . Some of these formulas contain the field  $S^*$ that is the sum of three vectors of curvature of the vector lines of the fields  $\tau, \nu, \beta$  and the field  $S_l$  that is the sum of three curvature vectors: vector lines of the normal fields  $\tau$  surfaces  $S_{\tau}$  and two vector lines of fields of their principle directions  $l_1$ ,  $l_2$ . In [8], it was found that the field  $S(\tau)$  is of the sum of three curvature vectors: the vector line  $L_{\tau}$  of the field  $\tau$  and any two geodesic lines (mutually orthogonal) on the surfaces  $S_{\tau}$ , orthogonal to the curves  $L_{\tau}$ . In the case when the curves  $L_{\tau}$  and the surface  $S_{\tau}$  are vector lines of the vector field  $\boldsymbol{v} = |\boldsymbol{v}|\boldsymbol{\tau}$  with the direction  $\boldsymbol{\tau}$  and the surfaces, orthogonal to them, the conservation laws obtained are equivalent to the divergence formulas for the field v.

With the help of these general geometric formulas, the differential conservation laws for the solutions of the eikonal  $|\operatorname{grad} \tau|^2 = n^2(x, y, z)$ , the Poisson equation  $\Delta u = -4\pi\rho(x, y, z)$  ( $\Delta u = u_{xx} + u_{yy} + u_{zz}$ ) and threedimensional solutions of Euler's hydrodynamic equations were obtained. In the plane case, the formulas found transform to conservation laws obtained in [2–4]. The symbols i, j, k denote the right-hand side system of unit vectors along the axes of the Cartesian coordinate system x, y, z. **Lemma 1** [2]. For any vector field  $\mathbf{v} = \mathbf{v}(x, y, z) = v_1 \mathbf{i} + v_2 \mathbf{j} + v_3 \mathbf{k} = |\mathbf{v}|\boldsymbol{\tau}$ defined in *D*, with components  $v_k(x, y, z) \in C^1(D)$ , k = 1, 2, 3, the modulus  $|\mathbf{v}| \neq 0$  in *D* and the direction  $\boldsymbol{\tau} = \mathbf{v}/|\mathbf{v}|$  ( $|\boldsymbol{\tau}| \equiv 1$ ) the following identity holds:

$$T(v) = S(\tau), \tag{1}$$

where

$$\begin{aligned} \boldsymbol{S}(\boldsymbol{\tau}) &\stackrel{\text{def}}{=} \operatorname{rot} \boldsymbol{\tau} \times \boldsymbol{\tau} - \boldsymbol{\tau} \operatorname{div} \boldsymbol{\tau} = (\boldsymbol{\tau} \cdot \nabla) \boldsymbol{\tau} - \boldsymbol{\tau} \operatorname{div} \boldsymbol{\tau}, \\ \boldsymbol{T}(\boldsymbol{v}) &\stackrel{\text{def}}{=} \operatorname{grad} \ln |\boldsymbol{v}| + \{\operatorname{rot} \boldsymbol{v} \times \boldsymbol{v} - \boldsymbol{v} \operatorname{div} \boldsymbol{v}\} / |\boldsymbol{v}|^2. \end{aligned}$$
(2)

By the direct calculation one can prove

**Lemma 2.** Let  $\boldsymbol{\tau} = \boldsymbol{\tau}(x, y, z) = \cos \alpha_1 \boldsymbol{i} + \cos \alpha_2 \boldsymbol{j} + \cos \alpha_3 \boldsymbol{k}$  be the vector field of the unit vectors  $(|\boldsymbol{\tau}| \equiv 1)$  with the domain of definition D,  $\alpha_1$ ,  $\alpha_2$ ,  $\alpha_3$  are the direction angles between the vector  $\boldsymbol{\tau}$  and the axes x, y, z, respectively, and  $\boldsymbol{\tau}(x, y, z) \in C^1(D)$ . Then the field  $\boldsymbol{S}(\boldsymbol{\tau})$  of the form of (2) can be represented in any of the forms  $\boldsymbol{S}(\boldsymbol{\tau}) = \sum_{j=1}^3 \operatorname{grad} \cos \alpha_j \times (\boldsymbol{i}_j \times \boldsymbol{\tau}) =$  $\sum_{j=1}^3 \cos \alpha_j \operatorname{rot}(\boldsymbol{\tau} \times \boldsymbol{i}_j), \ \boldsymbol{S}(\boldsymbol{\tau}) = \boldsymbol{\Phi}_1 - \operatorname{rot} \boldsymbol{\Psi} = \boldsymbol{\Phi}_2 + \operatorname{rot} \boldsymbol{\Psi}, \text{ where } \boldsymbol{i}_1 = \boldsymbol{i},$  $\boldsymbol{i}_2 = \boldsymbol{j}, \ \boldsymbol{i}_3 = \boldsymbol{k},$ 

From Lemmas 1 and 2 follows

**Theorem 1.** Under the conditions of Lemma 2 in D the following equivalent divergent identities (conservation laws) for the field  $\boldsymbol{\tau}$  are valid:  $\operatorname{div}\{\boldsymbol{S}(\boldsymbol{\tau}) - \boldsymbol{\Phi}_i\} = 0, \ i = 1, 2.$  If  $\boldsymbol{\tau}$  is the direction of the vector field  $\boldsymbol{v} = |\boldsymbol{v}|\boldsymbol{\tau}$ , then under the conditions of Lemma 1 in D the following equivalent divergent identities for the field  $\boldsymbol{v}$  hold:  $\operatorname{div}\{\boldsymbol{T}(\boldsymbol{v}) - \boldsymbol{\Phi}_i(\boldsymbol{v})\} = 0, \ i = 1, 2,$ where  $\boldsymbol{\Phi}_i(\boldsymbol{v})$  is obtained from  $\boldsymbol{\Phi}_i$  by replacing  $\boldsymbol{\tau}$  by  $\{\boldsymbol{v}/|\boldsymbol{v}|\}$  and  $\cos \alpha_j$  by  $\{v_j/|\boldsymbol{v}|\}.$ 

Let us obtain divergent formulas (differential conservation laws), which appear to be of a higher order for the field of the unit vectors  $\boldsymbol{\tau}$  or for a family of the curves  $\{L_{\tau}\}$ , as compared to the identities of Theorem 1. Let  $\{L_{\tau}\}$  be a family of the curves  $L_{\tau}$  with the tangent unit vector  $\boldsymbol{\tau} = \boldsymbol{\tau}(x, y, z)$ continuously filling the domain D. Let:

(D) one and only one curve  $L_{\tau} \in \{L_{\tau}\}$  passes through each point  $(x, y, z) \in D;$ 

- (E) at each point (x, y, z) of any curve  $L_{\tau} \in \{L_{\tau}\}$  there is a (right) Frenet basis  $(\tau, \nu, \beta)$ , so that three mutually orthogonal vector fields  $\tau, \nu, \beta$ are defined in D;
- (F)  $\boldsymbol{\tau}(x, y, z) \in C^2(D).$

By the direct calculations one can prove

**Lemma 3.** Let the family  $\{L_{\tau}\}$  of the curves  $L_{\tau}$  with the Frenet unit vectors  $\tau$ ,  $\nu$ ,  $\beta$ , the first curvature k and the second curvature  $\varkappa$  satisfy the conditions (D)–(F) in D. Let the field  $S^*$  be the sum of the three curvature vectors:

$$S^* \stackrel{\text{def}}{=} K_{\tau} + K_{\nu} + K_{\beta} = (\tau \cdot \nabla)\tau + (\nu \cdot \nabla)\nu + (\beta \cdot \nabla)\beta$$
  
=  $\operatorname{rot} \tau \times \tau + \operatorname{rot} \nu \times \nu + \operatorname{rot} \beta \times \beta$   
=  $-\{\tau \operatorname{div} \tau + \nu \operatorname{div} \nu + \beta \operatorname{div} \beta\} = \{S(\tau) + S(\nu) + S(\beta)\}/2.$ 

Here  $\mathbf{K}_{\tau} = (\boldsymbol{\tau} \cdot \nabla)\boldsymbol{\tau} = \operatorname{rot} \boldsymbol{\tau} \times \boldsymbol{\tau} = k\boldsymbol{\nu}, \ \mathbf{K}_{\nu} = (\boldsymbol{\nu} \cdot \nabla)\boldsymbol{\nu} = \operatorname{rot} \boldsymbol{\nu} \times \boldsymbol{\nu}, \ \mathbf{K}_{\beta} = (\boldsymbol{\beta} \cdot \nabla)\boldsymbol{\beta} = \operatorname{rot} \boldsymbol{\beta} \times \boldsymbol{\beta} \text{ are curvature vectors of the vector lines } L_{\tau}, \ L_{\nu}, \ L_{\beta} \text{ of the fields } \boldsymbol{\tau}, \ \boldsymbol{\nu}, \ \boldsymbol{\beta}, \text{ respectively. Then } \mathbf{S}^* = \mathbf{S}(\boldsymbol{\tau}) + \boldsymbol{\tau} \times \mathbf{R}^* \text{ in } D, \text{ where } \mathbf{R}^* \stackrel{\text{def}}{=} \varkappa \boldsymbol{\tau} + k\boldsymbol{\beta} + \boldsymbol{\beta} \operatorname{div} \boldsymbol{\nu} - \boldsymbol{\nu} \operatorname{div} \boldsymbol{\beta} = \boldsymbol{\Phi} + \mathbf{S}^* \times \boldsymbol{\tau} = \varkappa \boldsymbol{\tau} + (\boldsymbol{\tau} \cdot \operatorname{rot} \boldsymbol{\nu})\boldsymbol{\nu} + (\boldsymbol{\tau} \cdot \operatorname{rot} \boldsymbol{\beta})\boldsymbol{\beta}, \ \boldsymbol{\Phi} \stackrel{\text{def}}{=} \varkappa \boldsymbol{\tau} + k\boldsymbol{\beta}. \text{ Here } \boldsymbol{\Phi} \text{ is the Darboux vector } [5].$ 

Lemma 3 results in

**Theorem 2.** Under the conditions of Lemma 3 the divergent identity (conservation law for a family of curves  $\{L_{\tau}\}$ ) holds in D:

$$\operatorname{div}\{\boldsymbol{\tau}\operatorname{div}\boldsymbol{S}^* - \varkappa \operatorname{rot}\boldsymbol{\tau} - k \operatorname{rot}\boldsymbol{\beta}\} = 0 \quad \Leftrightarrow \\ \operatorname{div}\left\{\frac{1}{2}\boldsymbol{\tau}\operatorname{div}\boldsymbol{S}(\boldsymbol{\tau}) - k\boldsymbol{\nu}(\boldsymbol{\nu}\cdot\operatorname{rot}\boldsymbol{\beta}) - k\boldsymbol{\beta}[(\boldsymbol{\beta}\cdot\operatorname{rot}\boldsymbol{\beta}) + \varkappa]\right\} = 0.$$

Everywhere here the expression in braces is equal to  $\operatorname{rot} \mathbf{R}^*$ ;  $\operatorname{div} \mathbf{S}(\boldsymbol{\tau}) = 2(\boldsymbol{\tau} \cdot \operatorname{rot} \mathbf{R}^*)$ ;  $\operatorname{div} \mathbf{S}^* = (1/2) \operatorname{div} \mathbf{S}(\boldsymbol{\tau}) + k(\boldsymbol{\tau} \cdot \operatorname{rot} \boldsymbol{\beta}) + \varkappa(\boldsymbol{\tau} \cdot \operatorname{rot} \boldsymbol{\tau})$ ; the fields  $\mathbf{S}(\boldsymbol{\tau}), \, \mathbf{S}^*, \, \mathbf{R}^*, \, k, \, \varkappa$  are defined in Lemmas 1, 3.

The expressions for the value div  $S(\tau)$ , the first curvature k and the second curvature  $\varkappa$  of the curves  $L_{\tau}$  in terms of the fields of the Frenet unit vectors  $\tau$ ,  $\nu$ ,  $\beta$  are given by

**Lemma 4.** Under the conditions of Theorem 2, the following identities hold in D:

$$k = (\boldsymbol{\beta} \cdot \operatorname{rot} \boldsymbol{\tau}), \quad \boldsymbol{\varkappa} = \{(\boldsymbol{\tau} \cdot \operatorname{rot} \boldsymbol{\tau}) - (\boldsymbol{\nu} \cdot \operatorname{rot} \boldsymbol{\nu}) - (\boldsymbol{\beta} \cdot \operatorname{rot} \boldsymbol{\beta})\}/2,$$

$$\operatorname{div} \boldsymbol{S}(\boldsymbol{\tau}) = 2\{\boldsymbol{\varkappa}[\boldsymbol{\varkappa} - (\boldsymbol{\tau} \cdot \operatorname{rot} \boldsymbol{\tau})] - (\boldsymbol{\tau} \cdot [\operatorname{rot} \boldsymbol{\nu} \times \operatorname{rot} \boldsymbol{\beta}])\} \\ = 2\{(\boldsymbol{\nu} \cdot \operatorname{rot} \boldsymbol{\beta})(\boldsymbol{\beta} \cdot \operatorname{rot} \boldsymbol{\nu}) + [A^2 - (\boldsymbol{\tau} \cdot \operatorname{rot} \boldsymbol{\tau})^2]/4\} \\ A^2 + B^2 = (\operatorname{div} \boldsymbol{\tau})^2 + 2\operatorname{div} \boldsymbol{S}(\boldsymbol{\tau}) + (\boldsymbol{\tau} \cdot \operatorname{rot} \boldsymbol{\tau})^2,$$

where  $A \stackrel{\text{def}}{=} (\boldsymbol{\nu} \cdot \operatorname{rot} \boldsymbol{\nu}) - (\boldsymbol{\beta} \cdot \operatorname{rot} \boldsymbol{\beta}), \ \boldsymbol{B} \stackrel{\text{def}}{=} (\boldsymbol{\beta} \cdot \operatorname{rot} \boldsymbol{\nu}) + (\boldsymbol{\nu} \cdot \operatorname{rot} \boldsymbol{\beta}).$ 

Let us find a divergence formula (a conservation law) for the surfaces  $S_{\tau}$  given by some general properties in terms of their geometric characteristics. Let  $\{S_{\tau}\}$  be a family of the surfaces  $S_{\tau}$  with the unit normal  $\tau$  continuously filling the domain D in the space of x, y, z. The principal direction on  $S_{\tau}$  will be represented by a unit vector  $\mathbf{l}_i$  (i = 1, 2) with the corresponding direction;  $\mathbf{l}_i$  is the unit tangent vector of the curvature lines  $L_i$  on  $S_{\tau}$  [5–7]. Let:

- (A) through each point  $(x, y, z) \in D$  there passes one and only one surface  $S_{\tau} \in \{S_{\tau}\};$
- (B) at each point  $(x, y, z) \in D$  there exists a system (the right-hand side one) of the mutually orthogonal unit vectors  $\boldsymbol{\tau}$ ,  $\boldsymbol{l}_1$ ,  $\boldsymbol{l}_2$ , where  $\boldsymbol{\tau}$  is the unit normal,  $\boldsymbol{l}_1$  and  $\boldsymbol{l}_2$  are the principal directions on the surface  $S_{\tau}$ , passing through this point. For this, it is sufficient that each surface  $S_{\tau} \in \{S_{\tau}\}$  be  $C^2$ -regular [7]. Thus, three mutually orthogonal vector fields of the unit vectors  $\boldsymbol{\tau}(x, y, z)$ ,  $\boldsymbol{l}_1(x, y, z)$ ,  $\boldsymbol{l}_2(x, y, z)$  are defined in D. Simultaneously  $\boldsymbol{\tau} = \boldsymbol{\tau}(x, y, z)$  is the tangent unit vector of the curves  $L_{\tau}$ , orthogonal to the family  $\{S_{\tau}\}$ , that is, the vector lines of the normal vector field  $\boldsymbol{\tau}$ ;

(C)  $\tau \in C^{n}(D)$  (below n = 1 or 2),  $l_{i} \in C^{1}(D)$ , i = 1, 2.

The geometric meaning of the quantities div  $\tau$ , div  $S(\tau)$  and the divergent representations of the mean and the Gaussian curvatures of the surfaces is proved in

**Theorem 3** [8]. At any point  $(x, y, z) \in D$  under the conditions (A)–(C) the mean curvature H for n = 1 and the Gaussian curvature K for n = 2of the surface  $S_{\tau}$  passing through this point are equal to the divergence (the sources density) of the vector fields  $\{-\tau/2\}$  and  $\{-S(\tau)/2\}$ , respectively, at this point:  $H = -\frac{1}{2} \operatorname{div} \tau$ ,  $K = -\frac{1}{2} \operatorname{div} S(\tau)$ .

The geometric meaning of the conservation law of Theorem 1 explains

**Corollary 1.** Under the conditions of Theorem 3, the Gaussian curvature K of the surfaces  $S_{\tau} \in \{S_{\tau}\}$  admits in D the divergent representation of  $K = -\frac{1}{2} \operatorname{div} \Phi_i$ . If the surfaces  $S_{\tau}$  are orthogonal to the vector lines  $L_{\tau}$ 

of the vector field  $\mathbf{v} = |\mathbf{v}|\boldsymbol{\tau}, |\mathbf{v}| \neq 0$  in D and  $\mathbf{v} \in C^2(D)$ , then  $K = -\frac{1}{2} \operatorname{div} \mathbf{T}(\mathbf{v}) = -\frac{1}{2} \operatorname{div} \mathbf{\Phi}(\mathbf{v})$ . Here the fields  $\mathbf{S}(\boldsymbol{\tau}), \mathbf{T}(\mathbf{v}), \mathbf{\Phi}_i, \mathbf{\Phi}(\mathbf{v})$  are defined in Lemmas 1, 2.

The connection between the characteristics  $l_1$ ,  $l_2$ ,  $k_1$ ,  $k_2$ , H, K of the surfaces  $S_{\tau} \in \{S_{\tau}\}$  and the characteristics  $\tau$ ,  $\nu$ ,  $\beta$ , k,  $\varkappa$  of the curves  $L_{\tau}$  orthogonal to  $S_{\tau}$  is given by

**Theorem 4.** Let the family  $\{S_{\tau}\}$  of the surfaces  $S_{\tau}$  with the unit normal  $\tau = \tau(x, y, z)$  satisfy the conditions (A)–(C) for n = 2 and the family  $\{L_{\tau}\}$  of the curves  $L_{\tau}$ , orthogonal to  $\{S_{\tau}\}$ , satisfy the conditions (D)–(F). Then at each point  $(x, y, z) \in D$ , the principal directions  $l_1$  and  $l_2$ , the principal curvatures  $k_1$  and  $k_2$ , the mean curvature H and the Gaussian curvature K of the surface  $S_{\tau}$  passing through this point are expressed in terms of the Frenet unit vectors  $\tau$ ,  $\nu$ ,  $\beta$ , the first curvature k and the second curvature  $\varkappa$  of the curves  $L_{\tau}$  by the formulas:

$$l_{1} = \cos \omega \,\boldsymbol{\nu} + \sin \omega \,\boldsymbol{\beta}, \quad l_{2} = -\sin \omega \,\boldsymbol{\nu} + \cos \omega \,\boldsymbol{\beta},$$
  

$$\operatorname{tg} 2\omega = -\frac{A}{B} \quad \Leftrightarrow \quad (\boldsymbol{l}_{1} \cdot \operatorname{rot} \boldsymbol{l}_{1}) = (\boldsymbol{l}_{2} \cdot \operatorname{rot} \boldsymbol{l}_{2}),$$
  

$$k_{1} = -\frac{1}{2} \left\{ \operatorname{div} \boldsymbol{\tau} \pm \sqrt{A^{2} + B^{2}} \right\} = -(\boldsymbol{l}_{2} \cdot \operatorname{rot} \boldsymbol{l}_{1}),$$
  

$$k_{2} = -\frac{1}{2} \left\{ \operatorname{div} \boldsymbol{\tau} \mp \sqrt{A^{2} + B^{2}} \right\} = (\boldsymbol{l}_{1} \cdot \operatorname{rot} \boldsymbol{l}_{2}),$$
  

$$\Rightarrow \quad K \stackrel{\text{def}}{=} k_{1}k_{2} = \frac{1}{4} \left\{ (\operatorname{div} \boldsymbol{\tau})^{2} - (A^{2} + B^{2}) \right\}$$
  

$$K = (\boldsymbol{\tau} \cdot [\operatorname{rot} \boldsymbol{\nu} \times \operatorname{rot} \boldsymbol{\beta}]) - \varkappa^{2} = -\left\{ (\boldsymbol{\nu} \cdot \operatorname{rot} \boldsymbol{\beta})(\boldsymbol{\beta} \cdot \operatorname{rot} \boldsymbol{\nu}) + \frac{1}{4} A^{2} \right\}$$
  

$$\Rightarrow \quad H \stackrel{\text{def}}{=} \frac{1}{2} (k_{1} + k_{2}) = -\frac{1}{2} \operatorname{div} \boldsymbol{\tau}, \qquad K = -\frac{1}{2} \operatorname{div} \boldsymbol{S}(\boldsymbol{\tau}),$$

where the quantities A, B, k,  $\varkappa$ ,  $S(\tau)$ , div  $S(\tau)$  are given by the formulas of Lemmas 1, 2, 4. Let us place the sign "plus" in front of the radical for  $k_1 < k_2$  and "minus" for  $k_1 > k_2$ . We have  $H^2 - K = A^2 + B^2 \ge 0 \Rightarrow$  $H^2 \ge K$ .

**Lemma 5.** Let the conditions of Theorem 4 be valid and the field  $S_l^*$  be the sum of the three curvature vectors:

$$S_l^* \stackrel{\text{def}}{=} K_\tau + K_1 + K_2 = (\tau \cdot \nabla)\tau + (l_1 \cdot \nabla)l_1 + (l_2 \cdot \nabla)l_2$$
  
= rot  $\tau \times \tau$  + rot  $l_1 \times l_1$  + rot  $l_2 \times l_2$   
=  $-\{\tau \operatorname{div} \tau + l_1 \operatorname{div} l_1 + l_2 \operatorname{div} l_2\} = \{S(\tau) + S(l_1) + S(l_2)\}/2$ 

Here  $\mathbf{K}_{\tau} = (\boldsymbol{\tau} \cdot \nabla)\boldsymbol{\tau} = \operatorname{rot} \boldsymbol{\tau} \times \boldsymbol{\tau} = k\boldsymbol{\nu}$  is the curvature vector of the vector line  $L_{\tau}$  of the normal field  $\boldsymbol{\tau}$  of the surfaces  $S_{\tau}$ ,  $\mathbf{K}_{i} = (\mathbf{l}_{i} \cdot \nabla)\mathbf{l}_{i} = \operatorname{rot} \mathbf{l}_{i} \times \mathbf{l}_{i}$ is the curvature vector of the curvature lines  $L_{i}$  on  $S_{\tau}$  (i = 1, 2). Then in D

$$\begin{split} \boldsymbol{S}_{l}^{*} &= \boldsymbol{S}^{*} + \boldsymbol{\tau} \times \operatorname{grad} \boldsymbol{w}, \quad \boldsymbol{S}_{l}^{*} &= \boldsymbol{S}(\boldsymbol{\tau}) + \boldsymbol{\tau} \times \boldsymbol{R}_{l}^{*} \quad \Rightarrow \\ \boldsymbol{R}_{l}^{*} &\stackrel{\text{def}}{=} \operatorname{grad} \boldsymbol{w} + \boldsymbol{R}^{*} &= \varkappa_{l} \boldsymbol{\tau} + k\boldsymbol{\beta} + \boldsymbol{S}_{l}^{*} \times \boldsymbol{\tau} \\ &= \varkappa_{l} \boldsymbol{\tau} + \operatorname{rot} \boldsymbol{\tau} - (\boldsymbol{l}_{1} \operatorname{div} \boldsymbol{l}_{2} - \boldsymbol{l}_{2} \operatorname{div} \boldsymbol{l}_{1}) \\ &= \varkappa_{l} \boldsymbol{\tau} + \boldsymbol{l}_{1}(\boldsymbol{\tau} \cdot \operatorname{rot} \boldsymbol{l}_{1}) + \boldsymbol{l}_{2}(\boldsymbol{\tau} \cdot \operatorname{rot} \boldsymbol{l}_{2}), \end{split}$$

where  $S^*$ ,  $R^*$ , w are defined in Theorems 2, 4, and

$$\varkappa_{l} \stackrel{\text{def}}{=} -\{(\boldsymbol{l}_{1} \cdot \operatorname{rot} \boldsymbol{l}_{1}) + (\boldsymbol{l}_{2} \cdot \operatorname{rot} \boldsymbol{l}_{2})\}/2 = -(\boldsymbol{l}_{i} \cdot \operatorname{rot} \boldsymbol{l}_{i}), \quad i = 1, 2.$$

The expression of the conservation law of Theorem 2 in terms of the characteristics of the surfaces  $S_{\tau}$  is given by

**Theorem 5.** With the conditions and notations of Lemma 5 and Theorem 4 for the family  $\{S_{\tau}\}$  of the surfaces  $S_{\tau}$  in D there holds the divergent identity (conservation law):

$$\operatorname{div} \{ K\boldsymbol{\tau} + k_2 (\boldsymbol{l}_2 \cdot \operatorname{rot} \boldsymbol{\tau}) \boldsymbol{l}_1 - k_1 (\boldsymbol{l}_1 \cdot \operatorname{rot} \boldsymbol{\tau}) \boldsymbol{l}_2 \} = 0$$
  

$$\Leftrightarrow \quad \operatorname{div} \{ K\boldsymbol{\tau} + (H + B/2) \boldsymbol{K}_{\tau} - A \operatorname{rot} \boldsymbol{\tau}/2 \} = 0$$
  

$$\Leftrightarrow \quad \operatorname{div} \{ -\boldsymbol{\tau} \operatorname{div} \boldsymbol{S}_l^* + (\boldsymbol{l}_1 \cdot \operatorname{rot} \boldsymbol{\tau}) \operatorname{rot} \boldsymbol{l}_1 + (\boldsymbol{l}_2 \cdot \operatorname{rot} \boldsymbol{\tau}) \operatorname{rot} \boldsymbol{l}_2 + \varkappa_l \operatorname{rot} \boldsymbol{\tau} \}.$$

Here, the expression in braces  $\{ \}$  in everywhere equal to  $\{-\operatorname{rot} \mathbf{R}_l^*\} = -\operatorname{rot} \mathbf{R}^*$ ;  $K = -(\boldsymbol{\tau} \cdot \operatorname{rot} \mathbf{R}_l^*)$ ; div  $\mathbf{S}_l^* = -K + (\mathbf{l}_1 \cdot \operatorname{rot} \boldsymbol{\tau})(\boldsymbol{\tau} \cdot \operatorname{rot} \mathbf{l}_1) + (\mathbf{l}_2 \cdot \operatorname{rot} \boldsymbol{\tau})(\boldsymbol{\tau} \cdot \operatorname{rot} \mathbf{l}_2)$ .

**Remark 1.** The Frenet unit vectors  $\boldsymbol{\nu}, \boldsymbol{\beta}$  and the first curvature k of the curves  $L_{\tau}$  can be expressed in terms of  $\boldsymbol{\tau} \colon \boldsymbol{\nu} = (\operatorname{rot} \boldsymbol{\tau} \times \boldsymbol{\tau})/k, \, \boldsymbol{\beta} = \boldsymbol{\tau} \times \boldsymbol{\nu}, \, k = |\operatorname{rot} \boldsymbol{\tau} \times \boldsymbol{\tau}|$ . Since by the formulas of Lemma 4 and Theorem 4 the quantities  $\boldsymbol{l}_1, \boldsymbol{l}_2, k_1, k_2, \boldsymbol{\varkappa}, \boldsymbol{S}(\boldsymbol{\tau}), H, K$  are expressed in terms of the orts  $\boldsymbol{\tau}, \, \boldsymbol{\nu}, \, \boldsymbol{\beta}$  of the curves  $L_{\tau}$ , then finally, all these quantities can be expressed only through the field  $\boldsymbol{\tau}$ . Therefore, all the formulas of Theorems 2, 4, 5 and Lemmas 3–5 can be only expressed through the field  $\boldsymbol{\tau}$  (the unit tangent vectors of the curves  $L_{\tau}$  or normals to  $S_{\tau}$ ).

Let us apply the general formulas obtained to the solutions of the mathematical physics equations. Theorem 1 implies

**Corollary 2.** Let  $\tau = \tau(x, y, z)$  be the solution of the eikonal equation  $\tau_x^2 + \tau_y^2 + \tau_z^2 = n^2(x, y, z)$  in D, the time field  $\tau \in C^3(D)$ , the refractive index  $n \in C^2(D)$ . Then in D the following conservation law holds:

$$\operatorname{div}\{\boldsymbol{T} - \boldsymbol{\Phi}_i(\tau)\} = 0, \qquad i = 1, 2,$$

where  $\mathbf{T} = \operatorname{grad} \ln n - \Delta \tau \operatorname{grad} \tau/n^2$ ,  $\Phi_i(\tau)$  are obtained from the values  $\Phi_i$ of Lemma 2 by substitution  $\cos \alpha_1 = \tau_x/n$ ,  $\cos \alpha_2 = \tau_y/n$ ,  $\cos \alpha_3 = \tau_z/n$ .

**Corollary 3.** Let u = u(x, y, z) be the solution of the Poisson equation  $\Delta u = -4\pi\rho(x, y, z)$  in D,  $|\operatorname{grad} u| \neq 0$  in D, the potential  $u \in C^3(D)$ , the density  $\rho \in C^1(D)$ . Then, in D the conservation law div{ $T - \Phi_i(u)$ } = 0 holds, where  $T = \operatorname{grad} \ln |\operatorname{grad} u| + 4\pi\rho \operatorname{grad} u/|\operatorname{grad} u|^2$ , the fields  $\Phi_i(u)$  are defined in Corollary 2 by replacing  $\tau$  by u.

**Corollary 4.** Let  $\mathbf{v} = \mathbf{v}(z, y, z) = v\mathbf{\tau}$  be the velocity in Euler's hydrodynamic equations  $\mathbf{v}_t + \operatorname{grad} v^2/2 - \mathbf{v} \times \operatorname{rot} \mathbf{v} = \mathbf{F} - \operatorname{grad} p/\rho$ , which can be written down in D as  $\mathbf{G} = -\mathbf{T}(\mathbf{v}) (= -\mathbf{S}(\mathbf{\tau}))$ , where  $\mathbf{G} \stackrel{\text{def}}{=} \{\mathbf{v}_t + \mathbf{v} \operatorname{div} \mathbf{v} + \operatorname{grad} p/\rho - \mathbf{F}\}/v^2$ ;  $v \stackrel{\text{def}}{=} |\mathbf{v}| \neq 0$  in D,  $\mathbf{v} \in C^2(D)$ , the pressure  $p \in C^2(D)$ , the density  $\rho \in C^1(D)$ , the body force per unit of mass  $\mathbf{F} \in C^1(D)$ . Then in D the conservation law  $\operatorname{div}\{\mathbf{G} + \mathbf{\Phi}_i(\mathbf{v})\} = 0$ , i = 1, 2, holds, where the field  $\mathbf{\Phi}_i(\mathbf{v})$  is defined in Theorem 1.

Similarly by virtue of Theorems 2, 5, Remark 1, and the equality  $\boldsymbol{\tau} = \boldsymbol{v}/|\boldsymbol{v}|$ , we obtain for the above equations for the conservations laws of a higher order. In this case, the role of mutually orthogonal families of the curves  $L_{\tau}$  and the surfaces  $S_{\tau}$  for the eikonal equation is played by rays (the vector lines of the field  $\boldsymbol{v} = \operatorname{grad} \tau$ ) and by the fronts  $\tau(x, y, z) = \operatorname{const}$ ; for the Poisson equation—by the vector (force) lines of the field  $\boldsymbol{v} = \operatorname{grad} \boldsymbol{u}$  and by the equipotential surfaces  $u(x, y, z) = \operatorname{const}$ ; for Euler's hydrodynamic equations—by the streamlines (the vector lines of the velocity field  $\boldsymbol{v}$  at  $t = \operatorname{const}$ ) and by surfaces orthogonal to them. In the plane case, when all the quantities are independent of z, from Theorems 1, 2, 5 and from the Corollaries 2, 4 follow the conservation laws div  $\boldsymbol{S}(\boldsymbol{\tau}) = \operatorname{div} \boldsymbol{S}^* = 0$ , div  $\boldsymbol{T} = 0$ , and div  $\boldsymbol{G} = 0$  obtained in [2–4].

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